Heat & Mass Flow Processes

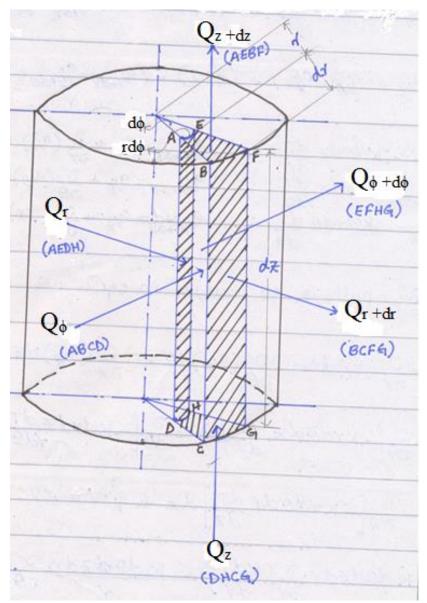
Week_03

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Let us consider a small element cut out of a pipe as shown in Fig 2.2. as given below in which heat enters from three faces and leaves the object from the other three faces in r, ϕ and z directions respectively.

Coordinates:	r	:	φ	:	Ζ
Heat Transfer:	Qr	•	Qø	•	Qz
Heat Flux (q=Q/A)	qr	:	q ø	•	qz
Thickness:	dr	:	rdø	:	dz



Energy balance for the small element is obtained from the First law of Thermodynamics.

(Net heat conducted into the element in dr, $rd\phi$, dz per unit time I) + (Internal heat generated per unit time II)

- = (Increase in internal energy per unit time III) + (work done by element per unit time IV)
- The IV term is small because the work done by the element due to change in temperature in neglected.

Qin–Qout + Qgen. = Rate of change in Internal Energy

$$(Q_r + Q_{\phi} + Q_z) - (Q_{r+dr} + Q_{\phi+d\phi} + Q_{z+dz}) + (Qgen.) = mc \frac{\partial T}{\partial \tau}$$
$$(Q_r + Q_{\phi} + Q_z) - (Q_{r+dr} + Q_{\phi+d\phi} + Q_{z+dz}) + (Qgen.) = \rho vc \frac{\partial T}{\partial \tau} \rightarrow (1)$$

Let

- ▷ Q_r be the heat flux in "r" direction at face AEDH and Q_{r+dr} be the heat flux at r+dr direction at face BCFG.
- ► Q_{ϕ} be the heat flux in " ϕ " direction at face ABCD and $q_{\phi+d\phi}$ be the heat flux at $\phi+d\phi$ direction at face EFHG.
- Q_z be the heat flux in "z" direction at face DHCG and Q_{z+dz} be the heat flux at z+dz direction at face AEBF

Now the rate of heat transfer in 'r', ' ϕ ', and 'z' direction is given by:

$$Q_{r} = q_{r}rd_{\phi}d_{z} = -k \ rd_{\phi}d_{z}\frac{\partial T}{\partial r}$$

$$Q_{\phi} = q_{\phi}d_{r}d_{z} = -k \ d_{r}d_{z}\frac{\partial T}{r\partial \phi}$$

$$Q_{z} = q_{z}rd_{r}d_{\phi} = -k \ rd_{r}d_{\phi}\frac{\partial T}{\partial z}$$

$$(2)$$

Similarly, the rate of heat transfer in 'r+dr', ' ϕ +d ϕ ', and 'z+dz' direction is given by:

$$Q_{r+dr} = Q_r + \frac{\partial Q}{\partial r} d_r \dots$$

$$Q_{\phi+d\phi} = Q_{\phi} + \frac{\partial Q}{\partial \phi} d_{\phi} \dots$$

$$Q_{z+dz} = Q_z + \frac{\partial Q}{\partial z} d_z \dots$$

$$(3)$$

Now by putting Equation (3) in Equation (1) we have.

$$\begin{aligned} &-\frac{\partial}{\partial r}(Q_{r})dr - \frac{\partial}{\partial \phi}(Q_{\phi})d\phi - \frac{\partial}{\partial z}(Q_{z})dz + q^{'''}v = \rho vc\frac{\partial T}{\partial \tau} \\ &-\frac{\partial}{\partial r}(-krd_{\phi}d_{z}\frac{\partial T}{\partial r})dr - \frac{\partial}{r\partial \phi}(-kd_{r}d_{z}\frac{\partial T}{\partial \phi})d\phi - \frac{\partial}{\partial z}(-krd_{r}d_{\phi}\frac{\partial T}{\partial z})dz + q^{'''}v = \rho vc\frac{\partial T}{\partial \tau} \\ &\frac{\partial}{\partial r}(krd_{\phi}d_{z}\frac{\partial T}{\partial r})dr + \frac{\partial}{\partial \phi}(kd_{r}d_{z}\frac{\partial T}{r\partial \phi})d\phi + \frac{\partial}{\partial z}(krd_{r}d_{\phi}\frac{\partial T}{\partial z})dz + q^{'''}v = \rho vc\frac{\partial T}{\partial \tau} \\ &kd_{r}d_{\phi}d_{z}\frac{\partial}{\partial r}(\frac{r\partial T}{\partial r}) + \frac{kd_{r}d_{\phi}d_{z}}{r}\frac{\partial}{\partial \phi}(\frac{\partial T}{\partial \phi}) + krd_{r}d_{\phi}d_{z}\frac{\partial}{\partial z}(\frac{\partial T}{\partial z}) + q^{'''}v = \rho vc\frac{\partial T}{\partial \tau} \end{aligned}$$

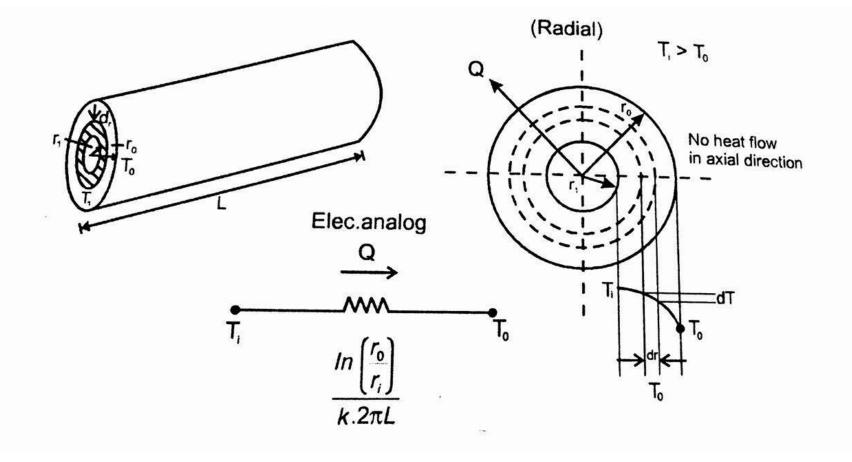
Since volume=v=dr.rd ϕ .dz, hence by dividing both sides of the above equation by "k.V", it is obtained that;

$$\frac{1}{r} \cdot \frac{\partial}{\partial r} \left[\frac{r \partial T}{\partial r} \right] + \frac{1}{r^2} \cdot \frac{\partial}{\partial \phi} \left[\frac{\partial T}{\partial \phi} \right] + \frac{\partial}{\partial z} \left[\frac{\partial T}{\partial z} \right] + \frac{q^{\prime \prime \prime}}{k} = \frac{1}{\alpha} \cdot \frac{\partial T}{\partial \tau} \rightarrow (4)$$

$$\left[\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial T}{\partial r} \right] + \left[\frac{1}{r^2} \cdot \frac{\partial^2 T}{\partial \phi^2} \right] + \left[\frac{\partial^2 T}{\partial z^2} \right] + \frac{q^{\prime \prime \prime \prime}}{K} = \frac{1}{\alpha} \cdot \frac{\partial T}{\partial \tau} \rightarrow (5)$$

Radial Heat Conduction Through a Cylinder without Internal Heat Generation

Consider a long hollow cylinder of length "L", inside radius r_i and outside radius r_o , having uniform thermal conductivity (k).



Radial Heat Conduction Through a Cylinder without Internal Heat Generation Assumptions:

- 1. There is no internal generation in the cylinder.
- 2. The inside and outside temperatures T_i and T_o are constant.
- 3. It's a long cylinder. So the end losses are negligible.

From general heat conduction equation;

$$\left[\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial T}{\partial r}\right] + \left[\frac{1}{r^2} \cdot \frac{\partial^2 T}{\partial \phi^2}\right] + \left[\frac{\partial^2 T}{\partial z^2}\right] + \frac{q^{\prime\prime\prime}}{K} = \frac{1}{\alpha} \cdot \frac{\partial T}{\partial \tau} \to (6)$$

Since there is steady state conduction in radial direction only without internal heat generation, all the other terms will become equal to zero. Hence equation (6) will become;

$$\left[\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial T}{\partial r}\right] = 0 \Longrightarrow \left[\frac{d^2 T}{dr^2} + \frac{1}{r} \cdot \frac{dT}{dr}\right] \Longrightarrow \frac{1}{r} \cdot \frac{d}{dr} \left[r \cdot \frac{dT}{dr}\right] = 0 \longrightarrow (7)$$

Heat Conduction Through a Cylinder without Internal Heat Generation

Since

$$\frac{1}{r} \neq 0 \Longrightarrow \frac{d}{dr} \left[r \cdot \frac{dT}{dr} \right] = 0 \to (8)$$

By double integrating equation (8) w.r.t. r;

$$\left[r.\frac{dT}{dr}\right] = C_1 \Rightarrow dT = C_1 r dr \Rightarrow T = C_1 \ln(r) + C_2 \rightarrow (9)$$

To find out C1 and C2, apply boundary conditions;

 $\begin{array}{ll} At & r=r_i & T=T_i \\ At & r=r_o & T=T_o \end{array} \end{array}$

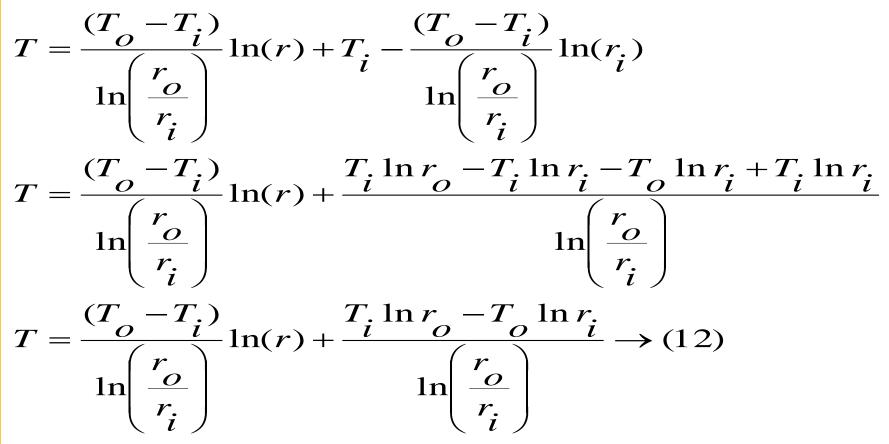
By applying boundary condition (1) and (2) in equation (9), we have;

$$T_i = C_1 \ln(r_i) + C_2 \rightarrow (10)$$
$$T_o = C_1 \ln(r_o) + C_2 \rightarrow (11)$$

<u>Heat Conduction Through a Cylinder without Internal Heat Generation</u> By solving equation (10) and (11), we have;

$$\begin{split} C_1 &= \frac{(T_i - T_o)}{\ln\left(\frac{r_i}{r_o}\right)} = \frac{(T_o - T_i)}{\ln\left(\frac{r_o}{r_i}\right)} \\ C_2 &= T_i - \frac{(T_o - T_i)}{\ln\left(\frac{r_o}{r_i}\right)} \ln(r_i) \\ & \ln\left(\frac{r_o}{r_i}\right) \end{split}$$

<u>Heat Conduction Through a Cylinder without Internal Heat Generation</u> Substituting C_1 and C_2 in equation (9)



This is called temperature distribution equation for cylinder system.

Heat Conduction Through a Cylinder without Internal Heat Generation The heat transfer rate in the cylinder can be determined by integrating the Fourier

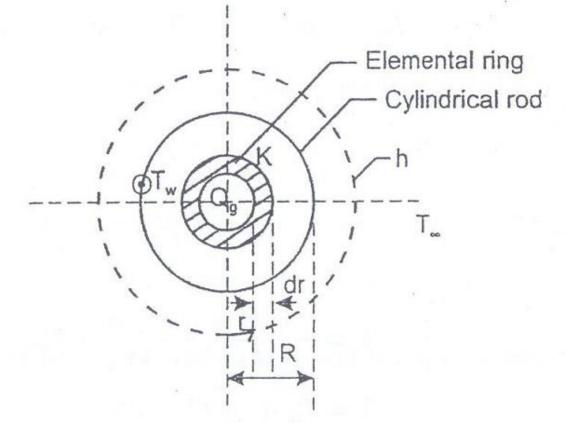
law of heat conduction as;

$$Q = -kA \frac{dT}{dr} \Rightarrow Qdr = -kAdT \Rightarrow Qdr = -k(2\pi.r.L)dT$$
$$\Rightarrow Q \int_{r_i}^{r_o} \frac{dr}{r} = -2\pi.k.L \int_{T_i}^{T_o} dT \Rightarrow Q \ln\left(\frac{r_o}{r_i}\right) = -2\pi kL(T_o - T_i)$$
$$\Rightarrow Q = \frac{2\pi kL(T_i - T_o)}{\ln\left(\frac{r_o}{r_i}\right)} = \frac{(T_i - T_o)}{\ln\left(\frac{r_o}{r_i}\right)/2\pi kL} = \frac{(T_i - T_o)}{R_{th}} \rightarrow (13)$$

Where R_{th} is the thermal resistance for hollow cylinder.

Radial Heat Conduction Through a Cylinder with Internal Heat Generation

Consider a cylinder of radius "R", having uniform thermal conductivity (k). Let an electric current passes through the cylinder, which causes uniform heat generation within the cylinder (q" W/m^3).ing of radius "r" thickness "dr" and heat generated in the elemental ring as Qg.



Radial Heat Conduction Through a Cylinder with Internal Heat Generation Assumptions:

- 1. There is internal generation in the cylinder.
- 2. The surface temperature and surrounding temperature are Tw and T α respectively.
- 3. It's a long cylinder. So the end losses are negligible.

From general heat conduction equation;

$$\left[\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial T}{\partial r}\right] + \left[\frac{1}{r^2} \cdot \frac{\partial^2 T}{\partial \phi^2}\right] + \left[\frac{\partial^2 T}{\partial z^2}\right] + \frac{q^{\prime\prime\prime}}{k} = \frac{1}{\alpha} \cdot \frac{\partial T}{\partial \tau} \to (14)$$

Since there is steady state conduction in radial direction only with internal heat generation, all the other terms will become equal to zero. Hence equation (14) will become;

$$\left[\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial T}{\partial r}\right] + \frac{q^{\prime\prime\prime}}{k} = 0 \rightarrow (15) \Rightarrow \left[\frac{d^2 T}{dr^2} + \frac{1}{r} \cdot \frac{dT}{dr}\right] + \frac{q^{\prime\prime\prime}}{k} = 0$$
$$\Rightarrow \left[r\frac{d^2 T}{dr^2} + \frac{dT}{dr}\right] + \frac{q^{\prime\prime\prime}}{k}r = 0 \Rightarrow \frac{dT}{dr}\left[r \cdot \frac{dT}{dr}\right] + \frac{q^{\prime\prime\prime}}{k}r = 0 \rightarrow (15)$$

<u>Heat Conduction Through a Cylinder with Internal Heat Generation</u> By double integrating equation (15), we have;

$$\begin{bmatrix} r \cdot \frac{dT}{dr} \end{bmatrix} + \frac{q^{\prime\prime\prime}}{k} \frac{r^2}{2} = C_1 \Rightarrow \begin{bmatrix} r \cdot \frac{dT}{dr} \end{bmatrix} = -\frac{q^{\prime\prime\prime}}{k} \frac{r^2}{2} + C_1 \Rightarrow \begin{bmatrix} \frac{dT}{dr} \end{bmatrix} = -\frac{q^{\prime\prime\prime}}{k} \frac{r}{2} + \frac{C_1}{r} \rightarrow (16)$$
$$T = -\frac{q^{\prime\prime\prime}}{k} \frac{r^2}{4} + C_1 \ln r + C_2 \rightarrow (17)$$

To find out C1 and C2, apply boundary conditions;

Atr = 0dT/dr = 0Atr = R $T = T_w$

By applying boundary condition (1) in equation (16), we have;

$$C_1 = 0$$

By applying boundary condition (2) in equation (17), we have;

$$T_{w} = -\frac{q'''}{k}\frac{R^{2}}{4} + (0)\ln r + C_{2} \Longrightarrow T_{w} = -\frac{q'''}{k}\frac{R^{2}}{4} + C_{2} \Longrightarrow C_{2} = T_{w} + \frac{q'''}{k}\frac{R^{2}}{4}$$

<u>Heat Conduction Through a Cylinder with Internal Heat Generation</u> By putting values of C1 and C2 in equation (17)

$$T = -\frac{q'''}{k}\frac{r^2}{4} + (0)\ln r + T_w + \frac{q'''R^2}{4k} \Longrightarrow T = T_w + \frac{q'''}{k}\left(R^2 - r^2\right) \to (18)$$

At center r=0, $T = T_{max}$

$$T_{\max} = T_w + \frac{q'''R^2}{k} \to (18)$$

Surface Temperature: Heat Generated inside the cylinder

$$q^{\prime\prime\prime} = \frac{Q}{A} \Longrightarrow Q = q^{\prime\prime\prime} A = q^{\prime\prime\prime} (\pi R^2 L) \to (19)$$

Heat due to convection on the surface of cylinder

$$Q = hA(T_w - T_\infty) = h(2\pi RL)(T_w - T_\infty) \to (20)$$

<u>Heat Conduction Through a Cylinder with Internal Heat Generation</u> Equating equation 19 and 20,

$$\pi R^2 Lq^{\prime\prime\prime} = h(2\pi RL)(T_w - T_\infty) \Longrightarrow Rq^{\prime\prime\prime} = 2h(T_w - T_\infty) \Longrightarrow 2hT_w = Rq^{\prime\prime\prime} + 2hT_\infty$$
$$T_w = T_\infty + \frac{Rq^{\prime\prime\prime}}{2h} \to (21)$$

Example 3.1:

A hollow cylinder 5 cm I.D and 10 cm O.D has an inner surface, temperature of $200^{\circ}C$ and an outer surface temperature of $100^{\circ}C$. Determine the temperature of the point half way between the inner and outer surfaces. If

the thermal conductivity of cylinder material is 70W/mK, determine also the heat flow through the cylinder per linear meter.

Solution

$$Q = \frac{-kA(T_0 - T_i)}{\ln\left(\frac{r_0}{r_i}\right)}$$
$$= \frac{kA(T_i - T_0)}{\ln\left(\frac{r_0}{r_i}\right)} = \frac{k \cdot 2\pi L(T_i - T_0)}{\ln\left(\frac{r_0}{r_i}\right)}$$

$$\frac{Q}{L} = \frac{70 \times 2 \times \pi \times (200 - 100)}{\ln \left(\frac{5}{2.5}\right)}$$
$$\frac{Q}{L} = 63434.3 \, W/m = 63.43 \, kW/m$$

At half-wave between r_i and $r_0 \therefore r' = \frac{r_i + r_0}{2} = \frac{5 + 2.5}{2} = 3.75 cm$

$$Q = \frac{2\pi k L(T_{i} - T_{0})}{\ln\left(\frac{r_{0}}{r_{i}}\right)} = \frac{2\pi k L(T_{i} - T)}{\ln\left(\frac{r'}{r_{i}}\right)}$$
$$(T_{i} - T) = \frac{(T_{i} - T_{0})}{\ln\left(\frac{r'}{r_{i}}\right)} = \frac{2\pi k L(T_{i} - T)}{\ln\left(\frac{r^{1}}{r_{i}}\right)}$$
$$(T_{i} - T) = \frac{(T_{1} - T_{0}) \ln\left(\frac{r'}{r_{i}}\right)}{\ln\left(\frac{r_{0}}{r_{i}}\right)}$$
$$= \frac{(200 - 100) \ln\left(\frac{3.75}{2.50}\right)}{\ln\left(\frac{5.00}{2.50}\right)} = 57.72$$
$$\therefore T = T_{i} - 57.72 = 200 - 57.72$$

 $T = 142.28^{\circ}C$

Example 3.2

A copper wire of 40 mm diameter carries 250 A and has a resistance of $0.25 \times 10^{-4} \Omega$ cm/length surface temperature of copper-wire is 250°C, and the ambient air temperature is 10°C. If the thermal conductivity of the copper wire is 175 W/mK. Calculate

- 1. Heat transfer co-efficient
- 2. Maximum temperature

Given:

$$D = 0.040 m$$

$$R = 0.020 m$$

$$I = 250 A$$

$$R_t = 0.25 \times 10^{-4} \Omega cm/length$$

$$T_w = 250 + 273 = 523 K$$

$$T_{\infty} = 10 + 273 = 283 K$$

$$k = 175 W/mK$$

Find:

(1) h (2) T_{\max}

Solution

We know that

Heat transfer $Q = I^2 R_t$ = $(250)^2 \times (0.25 \times 10^- 4)$ = 1.562 W/cm $Q = 1.562 \times 10^2 W/m$ Q = 156 W/m

We know that

Heat generated
$$q''' = \frac{Q}{V} = \frac{156}{\pi R^2 L}$$

= $\frac{156}{\pi (0.020)^2}$
 $q''' = 124140 W/m^3$

.:. Maximum Temperature

$$T_{\max} = T_w + \frac{q''' R^2}{4k}$$

= 523 + $\frac{124140 \times (0.020)^2}{4 \times 175}$
 $T_{\max} = 523.07 K$

I max *

To find heat transfer coefficient

$$T_{w} = T_{\infty} + \frac{q'''R}{2h}$$

523 = 283 + $\frac{0.020 \times 124140}{2 \times h}$
 $h = 5.17 W/m^{2\circ}C$

Links for Video Lectures

Theory Class:

https://pern.sharepoint.com/sites/HMFP-MECH-TECH-2018/Shared%20Documents/General/Recordings/HMFP%20THEORY-20210426_093607-Meeting%20Recording.mp4?web=1

Lab Session_03:

https://pern.sharepoint.com/sites/HMFP-MECH-TECH-2018/Shared%20Documents/General/Recordings/HMFP%20LAB-20210426_120950-Meeting%20Recording.mp4?web=1

